

## **CFD recombiner modelling and validation on the H2-Par and Kali-H2 experiments**

**S. Mimouni, N. Mechtoua**

*Electricité de France R&D Division  
6 Quai Watier F-78400 Chatou, France  
[stephane.mimouni@edf.fr](mailto:stephane.mimouni@edf.fr)*

**E. Moreau, M. Ouraou**  
*INCKA*

*85 av. Pierre Grenier F-92100 Boulogne, France*

### **ABSTRACT**

A large amount of Hydrogen gas is expected to be released within the dry containment of a pressurized water reactor (PWR), shortly after the hypothetical beginning of a severe accident leading to the melting of the core. According to local gas concentrations, the gaseous mixture of hydrogen, air and steam can reach the flammability limit, threatening the containment integrity. In order to prevent mechanical loads resulting from a possible conflagration of the gas mixture, French and German reactor containments are equipped with passive autocatalytic recombiners (PARs) which preventively oxidize hydrogen for concentrations lower than that of the flammability limit.

The objective of the paper is to present numerical assessments of the recombiner models implemented in CFD solvers NEPTUNE\_CFD and *Code\_Saturne*. NEPTUNE\_CFD is dedicated to the simulation of incompressible or compressible multi-component/multiphase flows. The multi-fluid set of equations is an extension of the “two fluid-one pressure” model to the case of  $m$  phases. Each fluid (fluid component and/or phase) is modelled through at least 3 conservation equations representing mass, momentum and total enthalpy. *Code\_Saturne* is dedicated to homogeneous incompressible or low Mach number compressible multi-component flows, with only one momentum equation, representing the momentum of the mixture of gas, liquid and particles. NEPTUNE\_CFD is mainly used for nuclear engineering, whereas *Code\_Saturne* is used for nuclear and fossil energy engineering, and for environment (geophysical flows).

Under the EDF/EPRI agreement, CEA has been committed to perform 42 tests of PARs. The experimental program named KALI-H2, consists checking the performance and behaviour of PAR. The CFD recombiner model implemented in the CFD codes is based on the manufacturer correlation for the hydrogen depletion rate. Concerning the KALI-H2 experiment, this model may lead to unrealistic values for the gas temperature, if the conjugate heat transfer and the wall steam condensation are not taken into account. The combined effects of these models give a good agreement between computational results and experimental data. Mesh sensitivity studies are systematically performed.

But if the recombiner design evolves, another manufacturer correlation would be required. An alternative would be to develop a numerical model to describe the PAR functioning and then deduce global correlations.

## **1 INTRODUCTION**

During a design-basis accident (DBA) or a severe accident (SA) in a nuclear power plant, certain chemical reactions may produce hydrogen (hydrogen is produced from the oxidation of zirconium sheaths and structures of fuel elements during the phase of core degradation), in such a way that hydrogen and oxygen volumetric concentrations may exceed the lower flammability limits (LFLs). The hydrogen risk in a nuclear power plant may be defined as the risk of hydrogen combustion in the containment building which represents a threat to the integrity of the confinement due to pressure and temperature levels. The nuclear safety is based on the concept called “defence-in-depth” that consists of a hierarchical deployment of different levels of equipment and procedures to maintain the efficiency of the physical barriers (International Atomic Energy Agency, 1996).

The transport and distribution of hydrogen inside the containment or the different compartments are critical phenomena to determine the kinetic and the nature of combustion. The prediction of stratification phenomena and location of hydrogen pockets are essential to assess the hydrogen risk and then, to optimise the hydrogen hazard mitigation system.

Among the different safety systems for limiting the pressure increase during the course of the accident and the impact of possible combustion (deflagration), French and German PWR reactors have three types of mitigation means.

- Sprinkler systems: the injected water droplets cool the containment and lower the pressure by condensing steam on the droplets. They also promote mixing of gas and rapidly break possible stratifications of the lightest gases.
- The walls of the containment building and metal structures also play an important role from a thermal viewpoint. The walls, significantly cooler than the gas, condense the water vapor in the gas mixture and thus limit the pressure increase in the containment. Furthermore, the temperature difference between fluid and walls generates a convection loop, enhancing the mixing of gases of different density.
- The passive auto-catalytic recombiners (PAR): their role is to pro-actively oxidize hydrogen for preventing its accumulation in the containment. The catalytic recombiners initiate a controlled combustion, which is similar to a slow deflagration.

A catalytic recombiner consists of a vertical channel or stack equipped with a catalyst bed in the lower part. In case of an accident, the catalyst is in contact with the gas mixture of the containment. Hydrogen molecules coming into contact with the catalyst surface are reacted with oxygen to form steam. The heat of the reaction at the catalyst surface causes a buoyancy-induced flow accelerating the inflow rate and thereby feeding the catalyst with a large amount of hydrogen that ensures high efficiency of recombination. The natural air circulation ensures a continuous air supply to the catalytic recombiner. This effect is increased by the height of the chimney, the inlet free area, and the fast heat up of the catalytic plates.

Catalytic recombiners favour the chemical reaction  $H_2 + \frac{1}{2} O_2 \rightarrow H_2O$  by lowering the activation energy threshold so that the reaction takes place at low temperature and concentration. The active catalyst materials include the noble metals platinum or palladium. The cover of the housing at the top protects the catalyst against direct spray of water and aerosol deposition. A catalytic recombiner is considered as « passive » because such a device is self-starting and self-feeding, and requires no external energy. Catalytic recombiners can start up with hydrogen fraction equal to about 2 %

This paper focuses on numerical assessments of PAR's modeling implemented in CFD solver *Code\_Saturne* and CMFD solver NEPTUNE\_CFD (Brailard,1999), (Bachelierie, 2003), (Avakian, 1999), (Fineschi, 1996), (Deng, 2008), (Kudriakov, 2006). It is organized as follows. The first part describes briefly the NEPTUNE\_CFD code (Mimouni, 2009). Next, the Gas dynamic model implemented in *Code\_Saturne* is described (Archambeau, 2004). The second part presents the LP manufacturer's models of recombiners and their implementation in these codes. The third part concerns the numerical assessments upon the H2PAR and KALIH2 experiments and the discussion of the results.

## 2 TWO-PHASE FLOW MODEL USED IN NEPTUNE\_CFD

The solver belongs to the well-known class of pressure based methods. It is able to simulate multi-component multiphase flows by solving a set of three balance equations for each field (fluid component and/or phase) (Ishii, 1975). These fields can represent many kinds of multiphase flows: distinct physical components (e.g. gas, liquid and solid particles); thermodynamic phases of the same component (e.g.: liquid water and its vapour); distinct physical components, some of which split into different groups (e.g.: water and several groups of different diameter bubbles); different forms of the same physical components (e.g.: a continuous liquid field, a dispersed liquid field, a continuous vapour field, a dispersed vapour field). The solver is implemented in the NEPTUNE software environment (Guelfi, 2007), (Mimouni, 2008), which is based on a finite volume discretization,

together with a collocated arrangement for all variables. The data structure is totally face-based which allows the use of arbitrary shaped cells (tetraedra, hexaedra, prisms, pyramids ...) including non-conformal meshes.

The main interest of the numerical method is the so-called "volume fraction – pressure – energy cycle" that ensures mass and energy conservation and allows strong interface source term coupling (Mechitoua, 2003).

Mass balance equations, momentum balance equations and total enthalpy balance equations are solved for each phase. The gas turbulence is taken into account by the classical k- $\epsilon$  model. The droplet diameter evolution is calculated from a equation of transport on the density of drops. Additional equations are added to take into account the non-condensable gases (air and hydrogen). As concern the interfacial momentum transfer terms, the only force exerted on droplet is the drag force. Small droplets stick at the wall and large drop slide along the wall under the competition between the surface tension and the gravity force. As a consequence, the gas velocity near the wall does not tend to zero but to the droplets velocity because of the drag force. This is a major difference between single-phase and two-phase flow approach (Mimouni, 2010). As concern the heat and mass transfer between droplets and the wall, it is based on the balance of heat and mass transfer between a drop and the gas mixture surrounding the drop using the correlations of Frössling/Ranz-Marshall which are of widespread use.

The model of drop-wall interaction which was developed and implemented is written as a symmetric extension of the nucleate boiling model at the wall, and uses as a starting point the model of mass transfer in the core flow. To establish this model, we made the following assumptions:

- the drops which accumulate on the walls take a hemispherical form;
- there is no nucleate boiling inside the drops at the wall;
- the drops which impact the walls successively see a stage of cooling (resp. heating) and a stage of condensation (resp. evaporation);
- the droplets stick to the wall (no rebound), or slide along the wall.

The total heat flux exchanged between the wall and the flow is split into four terms:

- $\varphi_{C1}$  a single-phase flow convective heat flux at the fraction of the wall area unaffected by the presence of droplets (heat transfer between the gas and the wall);
- $\varphi_{C2}$  a single-phase flow convective heat flux at the fraction of the wall area affected by the presence of a liquid film (heat transfer between the liquid film and the wall);
- $\varphi_{Th}$  a single-phase flow heat flux to decrease (resp. increase) the droplet temperature and reach the wall temperature (resp. the saturation state) (heat transfer between the droplets and the wall);
- $\varphi_E$  a condensation (resp. vaporisation) heat flux.

Details can be found in (Mimouni, 2010). An extensive validation process has been achieved in (Mimouni, 2010) against the COPAIN experiment and mesh sensitivity has been found acceptable.

### 3 HOMOGENEOUS GAS DYNAMIC MODEL USED IN CODE\_SATURNE

The motion, the distribution of gases and heat transfer in containment enclosures can be described by the general momentum, partial masses and energy conservation equations.

The predominant physical phenomena driving the motion, the distribution and heat transfer of fluids within containment enclosures are the following:

Mixing and / or segregation of gas whose velocity, density and temperature are different

- - "Swelling" of containment: the compressibility of gas is taken into account, even if the flow velocities are low.
- - Laminar and controlled combustion of hydrogen in recombiners, in order to limit the concentration of this gas
- - Condensation of steam on cold structure surfaces, which has the main effect of limiting the pressure rise.

The general momentum, partial masses and energy conservation equations describing these phenomena can be simplified and stiffness due to the presence of physics having very different characteristic length and time scales can be removed or relaxed.

The used turbulence model for containment applications is the standard k-epsilon one, supplemented by wall log laws for taking into account the turbulent friction and gaseous heat transfer between the fluid and the surrounding structures.

### 3.1 Low mach number approximation

The flows are mainly low Mach number flows, whose motion is predominantly driven by free convection. A low Mach number model can be implemented in a pressure correction based solver usually used for incompressible or steady dilatible flows, as Code\_Saturne (Archambeau, 2004). A spatial filtering of acoustic waves leads to separate the static pressure P into a uniform time-dependant thermodynamic pressure  $P_{th}(t)$  and a mechanical pressure  $p(x,t)$  :

$$P = P_{th}(t) + p(x,t); \quad P_{th} \gg |p(x,t)| \quad (1)$$

The general motion conservation equation of the mixture

$$\frac{\partial \rho \underline{V}}{\partial t} + \text{div} [ \underline{V} \otimes \rho \underline{V} - \mu_{tot} (\nabla \underline{V} + {}^t \nabla \underline{V}) + (2/3 \mu_{tot} \text{div} \underline{V} + 2/3 \rho k + P) \underline{I} ] = \rho \underline{g} + \Gamma_{cond} \underline{V}_l$$

associated with the supplementary approximations concerning the mechanical pressure and taking into account of mean hydrostatic pressure

$$p(x,t) \equiv p(x,t) + 2/3 \mu_{tot} \text{div} \underline{V} + 2/3 \rho k + \rho_0 g z \quad \text{with} \quad \rho_0 \equiv \int \rho d\Omega / \Omega_0$$

then becomes:

$$\frac{\partial \rho \underline{V}}{\partial t} + \text{div} [ \underline{V} \otimes \rho \underline{V} - \mu_{tot} (\nabla \underline{V} + {}^t \nabla \underline{V}) ] + \nabla p = (\rho - \rho_0) \underline{g} + \Gamma_{cond} \underline{V} \quad (2)$$

where  $\rho$ ,  $\underline{V}$ ,  $\mu_{tot}$ ,  $p$ ,  $\rho_0$ ,  $\Gamma_{cond}$  stand respectively for the mixture density, the mixture velocity vector, the total dynamic viscosity (including the turbulent viscosity deduced from the k-ε turbulence model), the mechanical pressure, the averaged density, the gravity acceleration and the condensation sink term.

Thanks to the low Mach number approximation, the mechanical pressure is neglected for the computation of density, through the thermal equation of state:

$$\rho = \frac{P_{th}}{RT \sum_k \frac{Y_k}{M_k}} \quad (3)$$

where R, T,  $Y_k$  and  $M_k$  stand respectively for the perfect gas Constant, the absolute temperature (in Kelvin), and the mass fractions of the different gases contained in the reactor building during a severe accident (Oxygen, Nitrogen, steam and Hydrogen).

The supplementary unknown  $P_{th}$  is solved, using integral forms over the entire domain  $\Omega_0$  of mass or enthalpy equations, written below.

$$\frac{\partial}{\partial t} \int_{\Omega_0} \rho d\Omega = \int_{\Omega_0} \Gamma_{cond} d\Omega - \int_{\partial\Omega_0} \rho \underline{V} \underline{n} dS \quad \text{or} \quad (4)$$

$$\frac{dP_{th}}{dt} \Omega_0 = - \frac{\partial}{\partial t} \int_{\Omega_0} \rho h d\Omega - \int_{\partial\Omega_0} [ \rho \underline{V} h - (\frac{\mu_t}{\sigma_t} + \frac{\lambda}{C_p}) \nabla h ] \underline{n} dS + \int_{\Omega_0} \rho S_h d\Omega$$

### 3.2 Energy equation

The enthalpy equation of the mixture is quite complex and contains several terms. The body forces, the viscous constraint contributions, the supplementary terms due to the presence of than more two different species are negligible, when compared to the convective and turbulent transport contributions. For low Mach number flows, the kinetic energy remains small when compared to the thermal energy. On the other side, the unsteady contribution of the thermodynamic pressure is conserved, as it plays a key role in the pressure rise in the containment.

The Fourier laminar and turbulent conduction term is directly written according to the enthalpy variable through the linearized relation:

$\lambda \frac{\partial T}{\partial x_j} = \frac{\lambda}{C_p} \frac{\partial h}{\partial x_j}$  where  $\lambda$  and  $C_p$  stand respectively for conductivity and specific Heat of the mixture.

The enthalpy equation is written in the following form:

$$\frac{\partial \rho h}{\partial t} + \text{div}[\rho \underline{V} h - (\frac{\mu_t}{\sigma_t} + \frac{\lambda}{C_p}) \nabla h] = \frac{dP_{th}}{dt} + \Gamma_{cond} h_{steam} - \rho \cdot E \cdot \frac{dm_{H2}}{dt} \quad (5)$$

where  $h$  stands for the sensitive enthalpy, defined as  $h \equiv \int_{T_0}^T C_p dT$ ,  $C_p$  is the mixture specific heat.

Then, in presence of exothermic chemical reactions, due to the combustion of hydrogen by the recombiners, the transformation of formation enthalpy into sensitive enthalpy is taken into account through a source term proportional to the sensitive enthalpy  $E$  released by the chemical reactions and to the chemical reaction speed  $\frac{dm_{H2}}{dt}$ . The same term is implemented in the energy equation in the

NEPTUNE\_CFD code.

We recall that the formation enthalpies and reaction heat  $E$  at ambient temperature are :

$$h_{H2}^0(T_0) = h_{He}^0(T_0) = h_{O2}^0(T_0) = h_{N2}^0(T_0) = 0; \quad h_{H2O,vap}^0(T_0) = -13.4 \text{ MJ/kg} \quad E = 1.22 \text{ MJ/kg}$$

The heat transfer due to condensation at the walls is modelled through a sink term proportional to the steam mass reduced into liquid water  $\Gamma_{cond}$  and the latent heat  $L_{cond}$ .

In the following sections, calculations performed with Code\_Saturne do not take into account the condensation phenomena at wall.

#### 4 RECOMBINERS MODELLING IN BOTH CODES

In this study we will use the following rate for the reaction  $H_2 + 1/2 O_2 \rightarrow H_2O$  :

$$\frac{dm_{H2}}{dt} = -\eta \cdot X_{min} \cdot (AP + B) \cdot \tanh[X_{H2} - 0.005] \quad \text{if} \quad X_{min} > 0.005 \quad (6)$$

$$\frac{dm_{H2}}{dt} = 0 \quad \text{if} \quad X_{min} < 0.005$$

Where  $dm_{H2}/dt$  stands for the recombined mass rate of  $H_2$  (kg/s), and where  $X_{min}$  is such as

$$X_{min} = \min(X_{H2}; 2X_{O2}; 0.08); \quad (7)$$

with  $X$  the molar fraction. Otherwise,  $P$  is the pressure in bars, and  $\eta$  a parameter interpreted as the recombiner output, and allowing to take into account the decrease of the efficiency of the recombiner for the weak concentrations of oxygen:

$$\begin{cases} \eta = 1 & \text{if } X_{H2} < X_{O2} \\ \eta = 0.6 & \text{if } X_{H2} > X_{O2} \end{cases} \quad (8)$$

The previous constants, given by SIEMENS for the FR90/1-150 PAR, take respectively the values  $0.48 \cdot 10^{-3}$  kg/s/bar and  $0.58 \cdot 10^{-3}$  kg/s (Braillard, 1999).

A important point to underline is that the hydrogen depletion rate is a semi-empirical relation, based on experimental measures made at the input and the output of the recombiner. In other words, the rate given by this expression is a global rate on the whole recombiner.

The reaction  $H_2 + 1/2 O_2 \rightarrow H_2O$  is exothermic. The released heat is equivalent to 1,22.108 J per kg of  $H_2$  burnt. The energy produced is released in the vicinity of the plates only if the reaction occurs, and is a function of the rate of  $H_2$  burnt. In the following,  $H_{reac}$ , is the enthalpy of the reaction (J/kg).

#### 5 MASS CONSERVATION EQUATION IN BOTH CODES

The mass conservation equations are written as below:

- the global mass equation, containing the sink term of wall condensation.

$$\frac{\partial \rho}{\partial t} + \text{div}(\rho \underline{V}) = \Gamma_{cond} \quad (9)$$

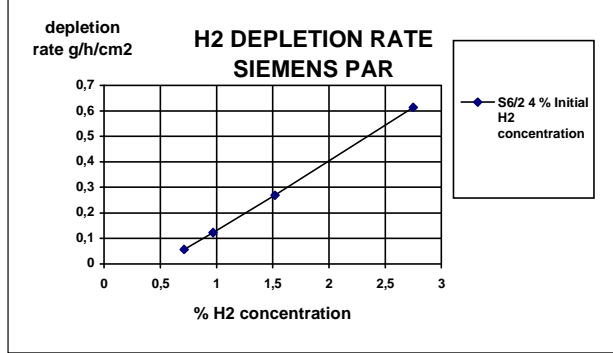


Figure 1: H2 depletion rate

- the conservation equations of non condensable gases, containing the slow combustion sink terms due to the recombiners:

$$\frac{\partial \rho_1 Y_{O_2}}{\partial t} + \text{div}(\rho_1 \underline{V} Y_{O_2} - \rho_1 D \nabla Y_{O_2}) = \Gamma_{O_2} = \rho_1 \frac{M_{O_2}}{2 M_{H_2}} \frac{dm_{H_2}}{dt}$$

$$\frac{\partial \rho_1 Y_{H_2}}{\partial t} + \text{div}(\rho_1 \underline{V} Y_{H_2} - \rho_1 D \nabla Y_{H_2}) = \Gamma_{H_2} = \rho_1 \frac{dm_{H_2}}{dt} \quad (10)$$

$$\frac{\partial \rho_1 Y_{N_2}}{\partial t} + \text{div}(\rho_1 \underline{V} Y_{N_2} - \rho_1 D \nabla Y_{N_2}) = \Gamma_{N_2} = 0$$

- the relation for obtaining the condensable gas (steam) from the concentration of the other gases.

$$Y_{H_2O} = 1 - Y_{O_2} - Y_{N_2} - Y_{H_2} \quad (11)$$

## 6 H2-PAR TEST CASE

To validate the model recombiners two H2PAR testing program were simulated and compared with experimental data. The experimental device consists of a sealed bag of flexible material to an approximate volume of 7.6 m<sup>3</sup>, which contains a Siemens FR90/1-150 type recombiners of 0.2m length, 0.166 m wide and 1.03 m in height (Note that the actual area of recombination corresponds to the Plate Size: 0.2 x 0.14mx 0.12m). The volume of this pocket is not constant over time, the variation being of the order of a m<sup>3</sup>. The atmosphere is initially composed of air and water vapor. Hydrogen is introduced at the base of the plant for a specified period. Hydrogen concentrations were measured by gas chromatography every minute somewhere in the area (input and output recombiners particular). They are then averaged over the whole field.

The computational domain is modelled by a cube and does not take into account the change in volume of the chamber (a flexible pouch) (Figure 2-Figure 3). The mesh contains 6859 cells, with 19 cells in each direction of space. The volume of recombiners is taken into account with approximately 40 stitches.

The operating conditions of test-E2 and E19 are summarized below:

### At inlet :

E2-bis test : The mass flowrate is successively  $1.3030 \cdot 10^{-3} \text{ kg.s}^{-1}$  (between 0 and 10s),  $4.3435 \cdot 10^{-4} \text{ kg.s}^{-1}$  (between 10 and 25s),  $1.9743 \cdot 10^{-4} \text{ kg.s}^{-1}$  (between 25 and 36s) . The hydrogen mass fraction is equal to 1 and the gas temperature is 85°C

E19 test : the mass flowrate is  $0.18 \cdot 10^{-3} \text{ kg.s}^{-1}$  between 0 and t 200s. The hydrogen mass fraction is 1 and the temperature is 70°C.

**Walls :** adiabatic conditions

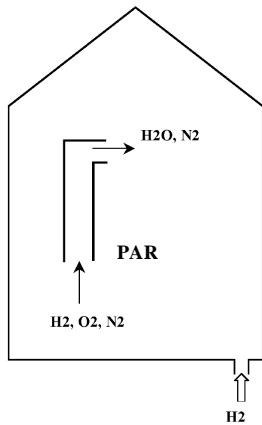


Figure 2: H2PAR test case

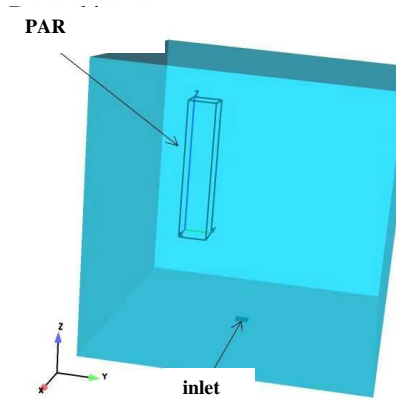


Figure 3: H2PAR test : computational domain

	Essai E2-Bis	Essai E19
O <sub>2</sub> Molar fraction	0.0828	0.1403
N <sub>2</sub> Molar fraction	0.3112	0.5274
H <sub>2</sub> O Molar fraction	0.6054	0.3323
H <sub>2</sub> Molar fraction	0	0
Temperature	85°C	70°C
Pressure	101300Pa	101300Pa

Table : initial conditions – H2PAR test

Figure 4 to Figure 7 show the velocity fields, temperature and hydrogen concentration on two planes of the mesh across the volume recombiners. The differences in density, generated by the combustion of hydrogen, create a natural convection loop that enhances the mixing at the same level and above recombiners. Moreover, cold stratified zone below the recombiners is calculated. The schematic modelling performed here may be representative of what can happen in the real case of the reactor building. Two opposing phenomena are in competition: the recombiners promote the controlled combustion of hydrogen, but according to their position in the enclosure, may also promote the creation of stratified layers of gas, "resistant" to the mix. Figure 8 and Figure 9 show the time evolution of the overall concentration of hydrogen dry gas (deduced from the volume fraction of water vapour) in the chamber and calculated with Code\_Saturne and scenario MAAP code with the 0D model. The continuous curves represent the results obtained when the source term is distributed over all the meshes representing the volume of the recombiner. The source term is calculated in each mesh recombiners with the local concentration of hydrogen and oxygen, weighted by the volume of the mesh on the overall volume of recombiner. The exercise carried out in this paragraph remains modest compared to the real complexity of the operation of a catalytic recombiner. But the goal here is to qualitatively reproduce the local effects induced by the operation of recombiners and more quantitative results to find the model-time when gas concentrations are homogeneous.

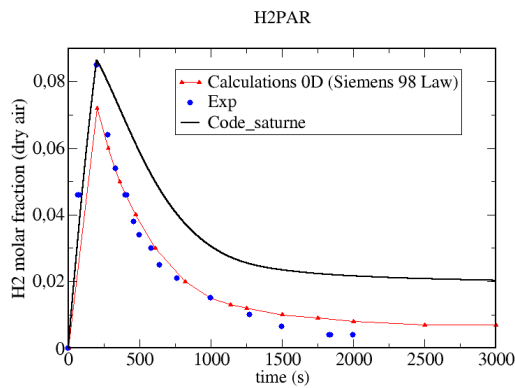


Figure 4: E19 test - Time evolution of hydrogen molar fraction

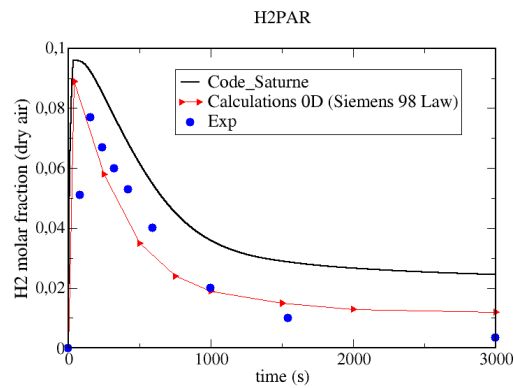


Figure 5: E2bis test – Time evolution of hydrogen molar fraction

## 7 KALI-H2 TEST

To test the catalytic capability of the SIEMENS recombiner model FR90 / 1-150, we used the KALI vessel. This 15.6 m<sup>3</sup> vessel (4.6m height and 2.1m of diameter) is able to withstand a pressure of 12 bar. The recombiner is located close to the wall, at the bottom of the vessel.

The vessel is connected with specific systems: steam injection system, hydrogen injection system, cold water system (Figure 10). Figure 11-Figure 13 show measurements location.

The vessel is also equipped with a fan to avoid any stratification during the hydrogen injection and also to start with a homogeneous mixture (it is stopped 45s after the beginning of the hydrogen injection). Since the test facility cannot accommodate a full size recombiner unit, a small size segment model was tested. Figure 12 shows that the small size recombiner is made of 15 cm height, 15 cm depth and 0.12 mm thickness with 1 cm spacing between them.

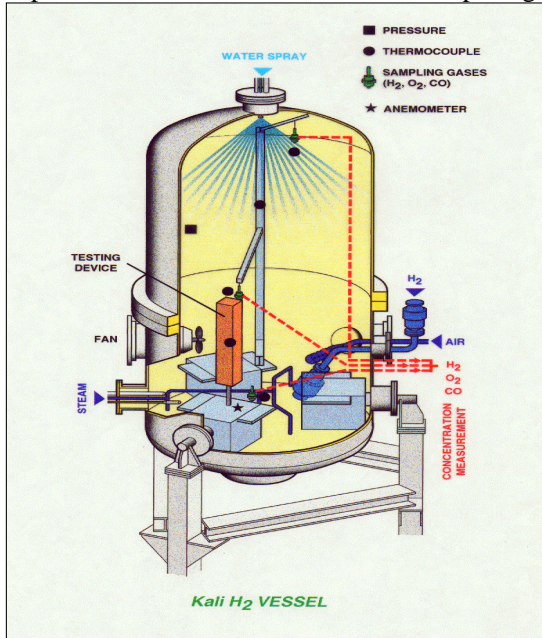


Figure 6: KALI-H2 test

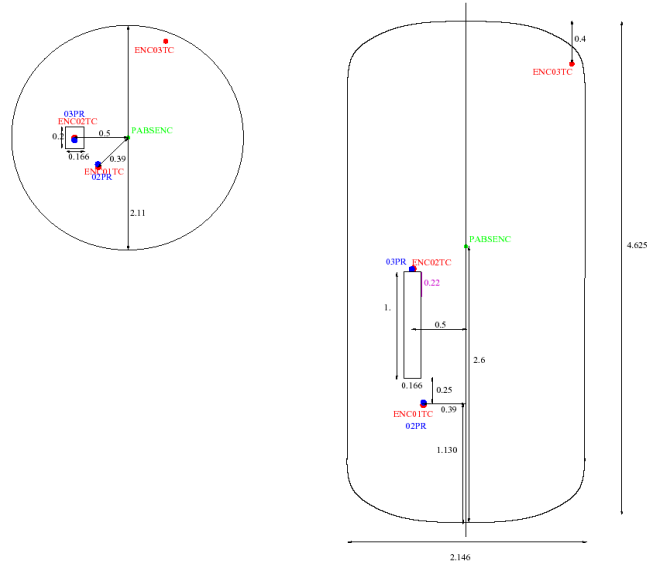
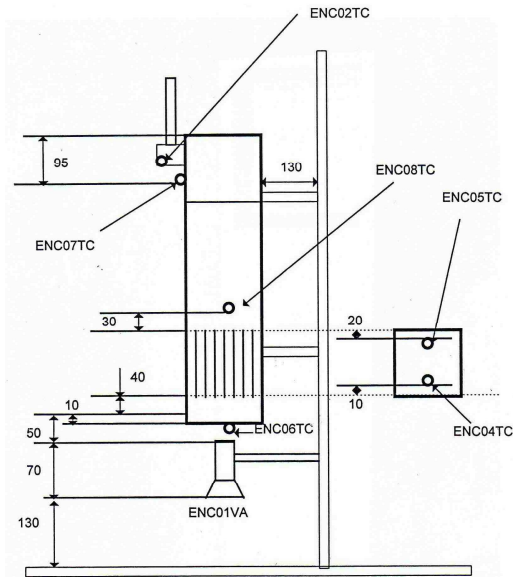
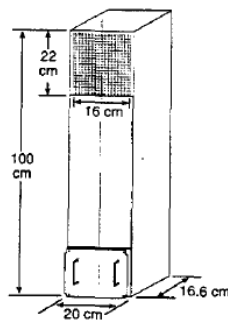


Figure 7: Sketch of KALI-H2 test



Only the plate located in the middle of the box is implemented with thermocouples

Figure 8: Sketch of KALI-H2 test



Siemens  
Segment Model  
PAR unit

Figure 9: Sketch of PAR unit



The recombiner was represented as a box of which the dimensions were the ones from the outside of the recombiner. Plates are not modeled.

In our global approach, source terms are distributed over the meshes representing the active volume of recombiner containing the vertical plates. Source terms are calculated in each mesh of the recombiner with the local concentration of hydrogen and oxygen, weighted by the active volume of the recombiner.

The simulations couple the two-phase flow and the solid heat conduction within the wall vessel. The wall condensation of vapor is taken into account with NEPTUNE\_CFD code. In fact, in a real situation, heat transfer to the wall and condensation result in an enhanced mixing of the atmosphere. In order to underline these aspects, Code\_Saturne simulations are performed with the artificial assumption of adiabatic walls and without wall condensation.

The initial conditions of the test 008 are the following:

- Pressure : 3.25 bars
- Temperature : 30°C (303 K)
- 4% of H<sub>2</sub>
- 96% of air.
- No steam (water) initially

The area is initially homogeneous for the mass fractions as well as for temperatures. The volume of the ½ domain is 7.42 m<sup>3</sup>. The total volume of H<sub>2</sub> is 4% \* 7.42 = 0.2968 m<sup>3</sup>. Thus, the total mass of H<sub>2</sub> is 0.077 kg ( $\rho \approx 0.26 \text{ kg/m}^3$ ). The total volume of air (H<sub>2</sub> and O<sub>2</sub>) is 96% \* 7.42 = 7.123 m<sup>3</sup>. Thus, the total mass of air is 26.5 kg ( $\rho \approx 3.72 \text{ kg/m}^3$ ).

As a consequence, the mass of non condensable gases is 26.58 kg.

Successive stages have to be considered for calculations : mixing of the gases with a fan device occurs every 900s for 120s. Moreover, an initial mixing is present during the first 120s of the test. The successive stages of the simulation are given in Figure 14.

We assume that the end of the mixing leads to an homogeneous gas mixing in the whole KALI-H2 vessel. Hence, the initial thermal-hydraulics properties of the calculations are taken at this time: homogeneous repartition of gas and temperature. The simulation carries on up to the second phase of mixing (900s → 1020s of test time). To take into account the homogenisation stage by the fan device, we reinitialise the computational values with averaged values calculated in the entire computational domain. This operation begins at the half of the duration of homogenisation process. Nevertheless, values inside the recombiner remain unchanged.

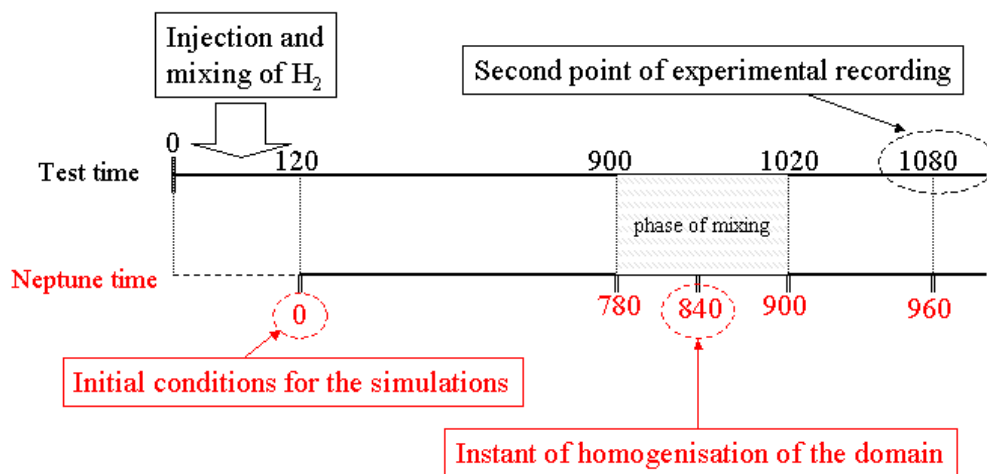


Figure 10: stages - test number 008

Then the mixing stage is simplified, but avoids taking into account the fan device for which few data are available. Then, calculations continue up to a second point of experimental recording (1080s).

The number of cells is about 112000. The time step corresponds to CFL=1. The sensitivity to the turbulence modelling has been investigated with NEPTUNE\_CFD code and calculations with three turbulent models (k-epsilon model, R<sub>ij</sub>-epsilon model (Mimouni, 2009, 2010), and laminar model) have been performed. The results are quite similar which means that the turbulence modelling does not affect the results significantly. In the following, the k-epsilon model is used in the calculations.

Figure 15 and Figure 16 represent the hydrogen mass fraction below the recombiner and at the outlet of the recombiner, calculated with NEPTUNE\_CFD and Code\_Saturne. Results are similar which means PAR models are implemented in a similar way in both codes. As a consequence, NEPTUNE\_CFD is indirectly checked against H2PAR test.

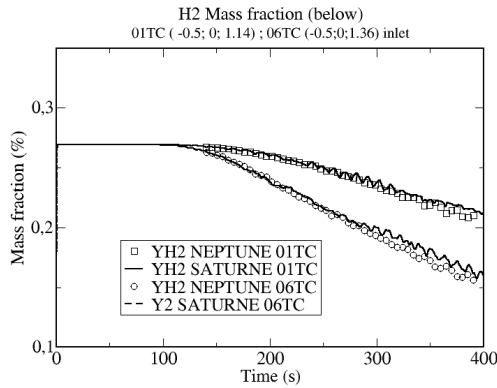


Figure 11: H2 Mass fraction below the PAR

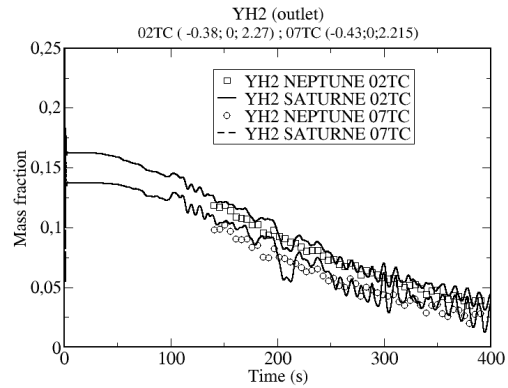


Figure 12: H2 Mass fraction at outlet of the PAR

Figure 17 shows the time evolution of gas temperature at the outlet of the PAR. We observe large discrepancies between Code\_Saturne and NEPTUNE\_CFD after 300s because the wall condensation effects and the heat exchange between the wall of vessel and the gas are not modelled in Code\_Saturne in this calculation. Firstly, gas temperature decreases because of conduction inside the wall (heat and mass exchange). Secondly, vapour condensates at the wall which means that pressure decreases inside the vessel and thus the gas temperature decreases too because of the gas perfect law. Calculations and experimental data are in reasonable agreement.

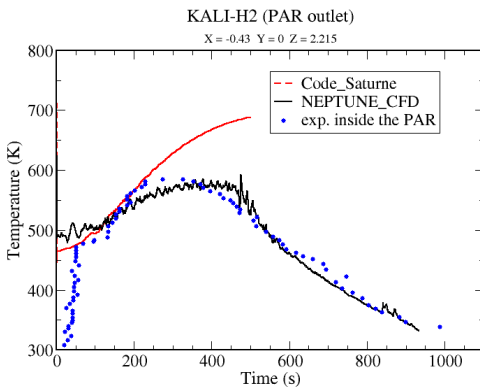


Figure 13: Time evolution of gas temperature at outlet of the PAR

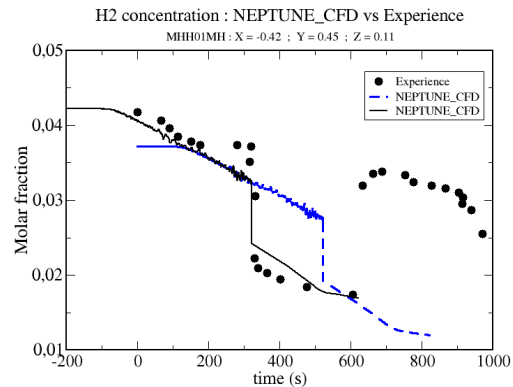


Figure 14: H2 Molar fraction just below the PAR

Figure 18 represents the hydrogen molar fraction just below the recombiner. The dashed line is the result directly calculated. The continuous line takes into account a start-up delay time described in (Avakian, 1999) and compares favourably with experimental values.

Figure 19 shows gas temperature and hydrogen mass fraction fields in the computational domain. The heat of the reaction at the catalyst surface causes a buoyancy-induced flow accelerating the inflow rate. Figure 19 clearly underlines the influence of the PAR outflow on the mixing process.

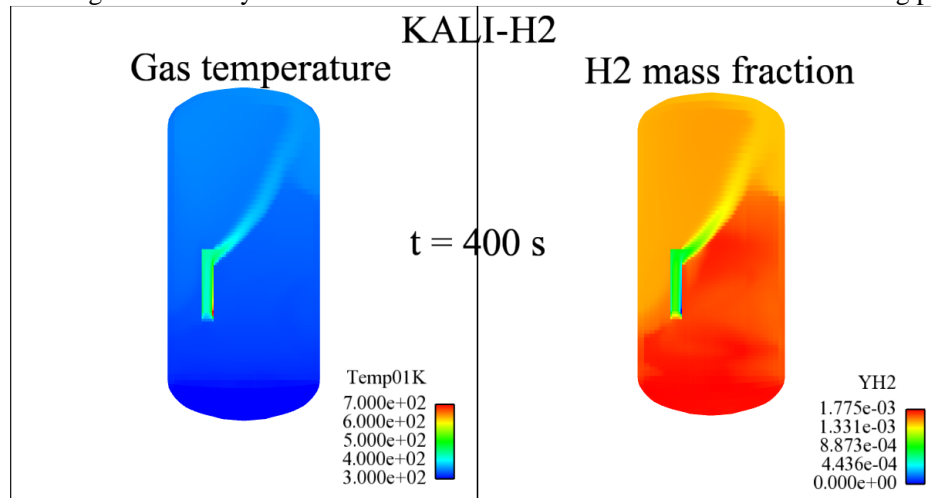


Figure 15: Gas temperature and hydrogen mass fraction fields.

## 8 CONCLUSION

We have presented in this paper the models implemented in NEPTUNE\_CFD, a three dimensional two-fluid code dedicated to nuclear reactor applications and in Code\_Saturne, a three dimensional single-phase code. Thanks to a code-to-experiment benchmark based on the COPAIN and TOSQAN facilities (Malet, 2008), we successfully evaluated the ability of the code to reproduce the vapor condensation at wall, atmosphere mixing and stratification in a vessel. These phenomena are of relevant interest in many industrial applications, especially regarding nuclear power plant containment at accident conditions. These models have been applied to the KALI-H2 test. Moreover, PAR models have been tested against H2PAR and KALI-H2 test and results are in a reasonable agreement with experimental values.

During the course of a severe accident in a Pressurized Water Reactor (PWR), spray systems are used in a containment in order to limit overpressure, to enhance the gas mixing in case of the presence of hydrogen and to drive down the fission products. Hence, vapor condensation on a cooled surface, spray effects and PAR systems act simultaneously in applications which is made possible with the two-phase flow approach proposed in the paper.

The PAR models used in the paper are based on the manufacturer's correlation to calculate the hydrogen depletion rate. In future works, it would be interesting to test the impact of the correlation uncertainties on gas temperature and hydrogen mass fraction at the PAR outlet calculated by a CFD code. Comparisons between manufacturer's correlation and more sophisticated models based on gas phase and surface chemical mechanisms should be investigated.

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